Can We Really Control Stimulated Raman and Brillouin Backscattering by Superposing Large Amplitude Waves onto the Plasma an Making it an Inhospitable Environment for Their Growth?

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Polymath Research Inc. $w_{n}^{2} = \frac{4\pi n e^{2}}{m}$ \mathcal{N}

The Big Picture

- Can we devise externally controllable methods of rendering the plasma conditions there where laser beams have to traverse in hohlraums (~ cm range) or direct drive ICF targets (~mm range) inhospitable to excessive levels of parametric instability and unchecked growth, backscattering, hot e⁻ generation, etc.?
- Could try and vary plasma density, temperature and velocity profiles by ab initio target fabrication and illumination condition designs (SSD, foam, PS). But this allows very little dynamical (run-time) control and have been shown to be weakly effective (< factor of 2).
- Alternative (or complementary) is to use optical mixing techniques, ie nonlinear optics technology to generate waves and disturbances which can do the job. How well? That is the subject of this study!

The Overall Program or Vision for OMC SSI Experiments





The Configuration of our Omega **Research Inc. Blue-Green OMC SSI Experiment**



SRS / SBS Saturation Wente 4-5-02

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Vlasov-Maxwell Simulations of Optical Mixing Generated EPWs (Nancy Collaboration) Blue-Green



OMC SSI Using IAWs: Experimental Configuration on Omega





Point the probe +/-500µm from target center along the target normal Point the pump +/-500µm from the taret center along target normal

What Are Our Goals with OMC SSI Omega Exp'ts?



- We measure how this <u>reduces the SRS and SBS</u> <u>backscattering levels</u> of the pump when the probe/pump energy ratio is high enough. See when effect starts and when it saturates per SRS wavelength or plasma density window near the sonic point.
- We hunt for laser beam localization or crossing volume localization effects by varying the pump focal position.

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2D Lasnex Simulations Suggest the Following Hydrodynamic Conditions for 10µm CH Foil Plasmas



Interaction beams start at 1.5 ns

2ns of SG3 DPP heaters 3kJ per ns (3 beams/ ns/ side). 2ns of SG3 DPP heaters 3kJ per ns (3 beams/ ns/ side).

Interaction beam starts at 1.5 ns

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Inferred (best fit) density profile has the shape:

 $n = n_{\text{peak}} \exp - \frac{z}{L_{\alpha}}^{\alpha}$

35%

time	n(peak)/nc	alpha	L_alpha microns	L_n microns	L_v microns	n(M=-1)/nc	Te keV peak	Te keV (M=-1)	Ti keV peak	A (M=-1)
t=1.5 ns	0.132	1.55	675	657	320	0.096	2.02	1.9	1.02	0.5
t=2.0 ns	0.085	1.5	867	811	440	0.059	2.08	2	0.97	0.23
t=2.5 ns	0.063	1.55	954	930	450	0.046	1.35	1.3	0.95	0.25

time	n(peak)/nc	alpha	L_alpha microns	L_n microns	L_v microns	n(M=-1)/nc	Te keV peak	Te keV (M=-1)	Ti keV peak	A (M=-1)
t=1.5 ns	0.146	1.55	667	636	343	0.102	2.07	1.92	1.03	0.55
t=2.0 ns	0.100	1.55	770	677	418	0.068	2.15	2.01	1.04	0.3
t=2.5 ns	0.065	1.6	1049	1061	431	0.051	1.53	1.47	0.95	0.27

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Transmission Spectra of the Probe When $I_{probe}/I_{pump} = 1/2$, Show Energy Transfer Especially at Late Time





Energy transfer Is Present Throughout Laser Pulse at $I_{probe}/I_{pump} = 1/15$ ie in the Small Signal Gain Regime







Probe Transmission Characteristics: Demonstration of IAW Generation



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Transmitted Beam Story Involves Euler's number?

100 ••••••1/2 T Probe (%) @ $1/2:1 T_{\text{Probe}}^{-}/T_{\text{Probe}}^{+}=1.21$ -1/2:0 T Probe (% 90 1/15 T Probe (%) @ $1/15:1 T_{Probe}|^{-} / T_{Probe}|^{+} = 2.6$ - 🖌 🛛 1.1 T Probe (%) 80 70 % T _{probe} 60 50 40 30 20 -500 0 500 Beam Location wrt Denisty Peak, µm

Polymath Research Inc. @1.1:1 $T_{\text{Probe}}^{-}/T_{\text{Probe}}^{+}=1.44$

@1.1:1 $T_{\text{Probe}}^{-}/T_{\text{pump off}}^{+}=1.13$ @1/2:1 $T_{\text{Probe}}^{-}/T_{\text{pump off}}^{+}=1.92$ @ $1/15:1 T_{Probe}^{-}/T_{pump off}^{+}=2.67$

> Note that @1.1:1 the transmission of the probe is LESS at +500 μ m than in the Pump off Case. This is because the Pump is now taking energy away from the Probe via **Optical Mixing as their roles have** reversed!

First Observation of Greater than 100% Crossed Beam Energy Transfer (Be Plasmas)



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Demonstration of SRS Suppression in the Appropriate Wavelength/density Window Dictated by the Localized IAW at the Mach -1 Surface









Comparison of SRS Power (GW) vs Time (ns) in 5 Wavelength Partitions at Mach -1



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Comparison of SRS Power (GW) vs Time (ns) in 5 Wavelength Partitions at Mach +1



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Comparison of SRS Power vs Time Mach +1 vs Mach -1 at 1:1





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Reassuringly, Very Little Difference Exists When Probe Intensity is Too Low (1/15:1) to Generate Larger Amplitude IAWs





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Reduction in SRBS is 4.8x in 500-515 nm Window when Large IAW Is Present



Interval in time where SRS at 1/15:1 looks different than SRS at 1:1 or 1/2:1 defines the time window used in this figure.

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Wavelength or Density Selectivity of SRS Suppression Is Due to Large Amplitude IAW Presence at Mach -1



Interval in time where SRS at 1/15:1 looks different than SRS at 1:1 or 1/2:1 defines the time window used in this figure.

Interval in space where SRS at 1/15:1 looks different than SRS at 1:1 or 1/2:1 corresponds to the wavelength range 490-515 nm.

This in turn corresponds to (assuming $T_e = 2 \text{ keV}$) the density range $0.04 < n_e / n_c < 0.052$ While 2D Hydro Simulations indicate $0.046 < n(M=-1) / n_c < 0.09$

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SRBS Reduction as a Function of Probe Energy Is Stronger in the Weak IAW Damping Limit

SRBS Energy (J)



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Conclusions



- We have demonstrated a strong suppression (x 5) in SRS backscattering reflectivity (normally of order 7%) in the presence of large amplitude ion acoustic waves in the strong IAW damping limit (10 μ m CH) and (x7) in the strong IAW damping limit (5 μ m Be).
- We have demonstrated that the wavelength range of SRS suppression overlaps strongly with the Mach -1 region of the plasma where a resonant IAW was driven by comparing to LASNEX predictions of hydro evolution.
- We have demonstrated that this should not be due to the seeding of LDI IAWs suppressing SRS, for even at 1/15:1 the seed source available for LDI would have been many orders of magnitude higher than thermal noise, yet +/-/0 focusing had no effect at 1/15:1 and all were like 0:1. Focusing changes only made a difference at higher energy ratios.
- We have seen > 100% transmission in the weak probe limit in Be targets making the energy transfer argument quite certain.
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Optical Diagnostics Used

- FABS stations deployed on BL 25 and BL30.
- Deploy Raman and Brillouin (3_0) channels on both.
- Streaked spectra and calorimetry data (time and wavelength averaged) on both SRS and SBS channels.
- Beam block calorimeter deployed on BL61 which is opposite BL30, which is our pump beam.
- This adds up to measuring: i) the Transmission of the Probe (BL 46) [streaked spectroscopy and calorimetry of BOTH SRS and SBS channels] at the FABS station on BL25,

ii) the reflectivity of the pump (BL30)

[streaked spectroscopy and calorimetry of BOTH SRS and SBS channels] at the FABS station on BL 30 and

iii) the transmission of the pump (BL30)

[just calorimetry of the SBS channel] at BL61.

• P510 data on all shots from BL46 and BL 30.